Interpolating the Derivatives of the Gamma Function

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Abstract

We consider a function $G(\lambda, z)$, entire in λ , which interpolates the derivatives of the Gamma function in the sense that $G(m, z) = \Gamma^{(m)}(z)$ for any integer $m \geq 0$ and we calculate the asymptotics of $G(\lambda, z)$ as *λ →* +*∞*.

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1 Introduction—the function $G(\lambda, z)$

Starting with the familiar integral representation of the Gamma function

$$
\Gamma(z) = \int_0^\infty e^{-t} t^{z-1} dt, \qquad \Re(z) > 0,
$$
\n(1)

and then taking derivatives, one obtains

$$
\Gamma^{(m)}(z) = \int_0^\infty e^{-t} (\ln t)^m t^{z-1} dt, \qquad m = 0, 1, 2, \dots, \quad \Re(z) > 0. \tag{2}
$$

The goal of this article is to construct and study an analytic function *G* of two variables λ and z which, for $\lambda = m$ becomes the *m*-th derivative $\Gamma^{(m)}(z)$ of the Gamma function.

To construct such a function we simply set

$$
G(\lambda, z) := \int_{C_r} e^{-\zeta} (\ln \zeta)^{\lambda} \zeta^{z-1} d\zeta, \qquad \Re(z) > 0,
$$
 (3)

where:

(i) The contour C_r consists of the intervals $[0, 1 - r]$ and $[1 + r, \infty)$ of the positive real semiaxis $(0 < r < 1)$ connected with the semicircle $\{\zeta \in \mathbb{C} : \zeta =$ $1 + re^{i\theta}$, $0 < \theta < \pi$, lying in the upper half-plane. The orientation of C_r is $0 \rightarrow \infty$.

(ii) $w = \ln \zeta$ denotes the branch of the logarithm which maps the open upper half-plane $\Im(\zeta) > 0$ conformally onto the open strip $0 < \Im(w) < \pi$. In particular, $\ln 1 = 0$ and $\ln(-1) = \pi i$.

(iii) $(\ln \zeta)^{\lambda} = e^{\lambda \ln \ln \zeta}$ and $\zeta^{z-1} = e^{(z-1)\ln \zeta}$. In particular $1^{\lambda} = 1^{z-1} = 1$.

It is clear that the value of the integral in (3) is independent of $r \in (0,1)$ and hence it defines a unique function $G(\lambda, z)$, which is entire in λ . Moreover, it is not hard to see that, as a function of λ , $G(\lambda, z)$ is of order 1 and of maximal (i.e. infinite) type for any *z* with $\Re(z) > 0$.

If $\Re(\lambda) > -1$ (and $\Re(z) > 0$), we can let $r \to 0^+$ in (3) and arrive at the formula

$$
G(\lambda, z) = e^{\pi i \lambda} \int_0^1 e^{-t} (-\ln t)^{\lambda} t^{z-1} dt + \int_1^{\infty} e^{-t} (\ln t)^{\lambda} t^{z-1} dt,
$$
 (4)

since for the negative values of ln *t* we have

$$
(\ln t)^{\lambda} = e^{\lambda \ln \ln t} = e^{\lambda [\ln(-\ln t) + \pi i]} = e^{\pi i \lambda} (-\ln t)^{\lambda}.
$$
 (5)

Formula (4) can be also written as

$$
G(\lambda, z) = \int_0^\infty e^{-t} (\ln t)^\lambda t^{z-1} dt, \qquad \Re(\lambda) > -1, \quad \Re(z) > 0,
$$
 (6)

from which we can see immediately that $G(m, z) = \Gamma^{(m)}(z)$ for $m = 0, 1, \ldots$ and $\Re(z) > 0$.

2 Some basic properties of $G(\lambda, z)$

Before we start our analysis let us introduce the notation

$$
\mathbb{C}_{\xi} := \{ z \in \mathbb{C} : \Re(z) > \xi \},\tag{7}
$$

where ξ is a real number.

As we have, essentially, already observed, it follows from (3) that $G(\lambda, z)$ is analytic in $\mathbb{C} \times \mathbb{C}_0$.

We continue with some basic properties of $G(\lambda, z)$.

Property 1. For any $\lambda \in \mathbb{C}_{-1}$ we have

$$
L(\lambda) := \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} G(\lambda, z) = e^{\pi i \lambda} \Gamma(\lambda + 1),
$$
\n(8)

where $z^{\lambda+1} = e^{(\lambda+1)\ln z}$, $\ln z$ being continuous in \mathbb{C}_0 (in particular $1^{\lambda+1} = 1$).

Proof. In view of (6) we have

$$
G(\lambda, z) = \int_0^{\varepsilon} e^{-t} (\ln t)^{\lambda} t^{z-1} dt + \int_{\varepsilon}^{\infty} e^{-t} (\ln t)^{\lambda} t^{z-1} dt, \quad \Re(\lambda) > -1, \ \Re(z) > 0,
$$
\n(9)

for any fixed $\varepsilon > 0$. Now, the integral

$$
\int_{\varepsilon}^{\infty} e^{-t} (\ln t)^{\lambda} t^{z-1} dt \tag{10}
$$

is entire in *z* and, in particular, finite at $z = 0$; thus it does not contribute to the value of the limit $L(\lambda)$. Consequently, for the limit in (8) we have

$$
L(\lambda) = \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} \int_0^{\varepsilon} e^{-t} (\ln t)^{\lambda} t^{z-1} dt,
$$
\n(11)

independently of ε , as long as $\varepsilon > 0$. Since for $t \in [0, \varepsilon]$ the quantity e^{-t} can be made as close to 1 as we wish by choosing ε sufficiently small, it follows that

$$
L(\lambda) = \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} \int_0^\varepsilon (\ln t)^{\lambda} t^{z - 1} dt.
$$
 (12)

But the limit in (12) is, also, independent of ε , as long as $\varepsilon > 0$. Therefore, we can take $\varepsilon = 1$ in (12) and get

$$
L(\lambda) = \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} \int_0^1 (\ln t)^{\lambda} t^{z-1} dt
$$

$$
= e^{\pi i \lambda} \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} \int_0^1 (-\ln t)^{\lambda} t^{z-1} dt
$$

$$
= e^{\pi i \lambda} \lim_{\substack{z \to 0 \\ z \in \mathbb{C}_0}} z^{\lambda + 1} \int_0^\infty e^{-z\tau} \tau^{\lambda} d\tau
$$
(13)

(the last integral is obtained by substituting $\tau = -\ln t$ in the previous integral). Finally, since for real and positive *z*, the substitution $t = z\tau$ yields

$$
z^{\lambda+1} \int_0^\infty e^{-z\tau} \tau^\lambda d\tau = \Gamma(\lambda+1), \qquad \lambda \in \mathbb{C}_{-1},
$$
 (14)

it follows by analytic continuation that the equality (14) holds for any $z \in \mathbb{C}_0$. Ξ

Property 2. If ∂_z denotes the derivative operator with respect to *z*, then

$$
\partial_z G(\lambda, z) = G(\lambda + 1, z), \qquad \lambda \in \mathbb{C}, \quad z \in \mathbb{C}_0.
$$
 (15)

Property 2 follows by differentiating (3) with respect to *z*.

Property 3. The function $G(\lambda, z)$ satisfies the functional equation

$$
G(\lambda, z + 1) = zG(\lambda, z) + \lambda G(\lambda - 1, z), \qquad \lambda \in \mathbb{C}, \quad z \in \mathbb{C}_0.
$$
 (16)

Proof. For $\Re(\lambda) > 1$ and $\Re(z) > 0$ integration by parts yields

$$
z \int_0^{\infty} e^{-t} (\ln t)^{\lambda} t^{z-1} dt = e^{-t} (\ln t)^{\lambda} t^{z} \Big|_{t=0}^{\infty} + \int_0^{\infty} \left[e^{-t} (\ln t)^{\lambda} t^{z} - \lambda e^{-t} (\ln t)^{\lambda - 1} t^{z-1} \right] dt
$$

$$
= \int_0^{\infty} e^{-t} (\ln t)^{\lambda} t^{z} dt - \lambda \int_0^{\infty} e^{-t} (\ln t)^{\lambda - 1} t^{z-1} dt
$$

$$
= G(\lambda, z + 1) - \lambda G(\lambda - 1, z).
$$

Thus, $G(\lambda, z)$ satisfies (16) for $\Re(\lambda) > 1$ and, consequently, for all $\lambda \in \mathbb{C}$ by analytic continuation.

We can now combine Property 2 with Property 3 and get

$$
G(\lambda, z + 1) = z \partial_z G(\lambda - 1, z) + \lambda G(\lambda - 1, z), \qquad \lambda \in \mathbb{C}, \quad z \in \mathbb{C}_0. \tag{17}
$$

Multiplying (17) by $z^{\lambda-1}$ yields

$$
z^{\lambda - 1} G(\lambda, z + 1) = \partial_z \left[z^{\lambda} G(\lambda - 1, z) \right], \qquad \lambda \in \mathbb{C}, \quad z \in \mathbb{C}_0, \tag{18}
$$

so that, by integrating (18) from z_0 to z we obtain

$$
z^{\lambda}G(\lambda - 1, z) - z_0^{\lambda}G(\lambda - 1, z_0) = \int_{z_0}^{z} \zeta^{\lambda - 1}G(\lambda, \zeta + 1) d\zeta, \qquad \lambda \in \mathbb{C}, \quad (19)
$$

where the integral is taken over a rectifiable arc joining z_0 and z , and lying entirely in the open right half-plane \mathbb{C}_0 .

Incidentally, if we restrict $\lambda \in \mathbb{C}_0$, then we can let $z_0 \to 0$ in (19) and invoke Property 1 in order to get

$$
z^{\lambda}G(\lambda - 1, z) + e^{\pi i \lambda} \Gamma(\lambda) = \int_0^z \zeta^{\lambda - 1} G(\lambda, \zeta + 1) d\zeta, \qquad \lambda, z \in \mathbb{C}_0.
$$
 (20)

Noticing that, for any $\lambda \in \mathbb{C}$, the integral in the right-hand side of (19) defines a (multivalued) analytic function of *z* in the half-plane C*−*¹ we can conclude that $z^{\lambda}G(\lambda-1, z)$ and, consequently, $G(\lambda-1, z)$ have a meromorphic extension for $z \in \mathbb{C}_{-1}$, with a branch point at $z = 0$. It is convenient (and harmless) to denote by $G(\lambda - 1, z)$ too the meromorphic extension of $G(\lambda - 1, z)$.

It, then, follows that for a fixed $z_0 \in \mathbb{C}_0$ formula (19) remains valid for $z \in \mathbb{C}_{-1}$ (as long as the integral is taken over an arc joining z_0 and z , and lying entirely in \mathbb{C}_{-1}). Therefore, by the same argument $G(\lambda - 1, z)$ has a (multivalued) meromorphic extension for $z \in \mathbb{C}_{-2}$. As the dummy variable ζ of the integral in (19) takes the value -1 , the function $G(\lambda, \zeta + 1)$ appearing in the integrand is evaluated at its branch point 0. Hence *z* = *−*1 is a branch point of the integral and, hence, of $G(\lambda - 1, z)$. Keeping arguing in the same manner, we can conclude that $G(\lambda - 1, z)$ has a (multivalued) meromorphic extension in \mathbb{C}_{-n} for any $n = 1, 2, \ldots$, with branch points at $z = 0, 1, \ldots, n - 1$. And since *n* is arbitrary, we can finally conclude that $G(\lambda - 1, z)$ has a (multivalued) meromorphic extension in $\mathbb C$ with branch points at the nonpositive integers. For the integral values of λ , namely for $\lambda = m$, $m = 1, 2, \dots$, we know that $G(\lambda - 1, z) = \Gamma^{(m-1)}(z)$ has poles of order *m* at $z = -n$, where $n = 0, 1, 2, \ldots$. From the above discussion it also follows that the limit in (8) of Proposition 1 remains valid as long as *z* approaches zero in a continuous fashion, without ever crossing the negative real semiaxis.

Let us summarize the above observations in the following property.

Property 4. For any complex λ , the function $G(\lambda, z)$, viewed as a function of z , has a (multivalued) meromorphic extension in $\mathbb C$ with branch points at $z = 0, -1, -2, \ldots$. Furthermore, Property 1 can be strengthen as

$$
\lim_{\substack{z \to 0 \\ z \in \gamma}} z^{\lambda + 1} G(\lambda, z) = e^{\pi i \lambda} \Gamma(\lambda + 1),\tag{21}
$$

where γ is a continuous curve which does not cross the negative real semiaxis, while $z^{\lambda+1}$ is defined so that it is continuous in *z* and $1^{\lambda+1} = 1$.

Finally, from (19) it follows easily the following property:

Property 5. For any $z \neq 0, -1, -2, \ldots$, the function $G(\lambda, z)$ is entire in λ .

3 Asymptotics of $G(\lambda, z)$ as $\lambda \to +\infty$

For typographical convenience we write (4) as

$$
G(\lambda, z) = G_0(\lambda, z) + G_1(\lambda, z), \tag{22}
$$

where

$$
G_0(\lambda, z) := e^{\pi i \lambda} \int_0^1 e^{-t} (-\ln t)^{\lambda} t^{z-1} dt, \qquad \Re(\lambda) > -1, \quad \Re(z) > 0, \tag{23}
$$

and

$$
G_1(\lambda, z) := \int_1^\infty e^{-(t-\lambda \ln \ln t)} t^{z-1} dt, \qquad \Re(\lambda) > -1, \quad \Re(z) > 0 \tag{24}
$$

(in fact, $G_1(\lambda, z)$ is entire in z).

3.1 The expansion of $G_0(\lambda, z)$

By expanding e^{-t} in (23) as a power series we obtain

$$
G_0(\lambda, z) = e^{\pi i \lambda} \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \int_0^1 (-\ln t)^{\lambda} t^{n+z-1} dt, \qquad \Re(\lambda) > -1, \quad \Re(z) > 0.
$$
\n(25)

Now, as in (13)-(14),

$$
\int_0^1 (-\ln t)^\lambda t^{n+z-1} dt = \int_0^\infty e^{-(n+z)\tau} \tau^\lambda d\tau = \frac{\Gamma(\lambda+1)}{(n+z)^{\lambda+1}}.\tag{26}
$$

Thus, (25) becomes

$$
G_0(\lambda, z) = e^{\pi i \lambda} \Gamma(\lambda + 1) \sum_{n=0}^{\infty} \frac{(-1)^n}{n! (n+z)^{\lambda+1}}, \qquad \Re(\lambda) > -1, \quad \Re(z) > 0, \quad (27)
$$

and the series converges very rapidly. Since for $\Re(z) > 0$

$$
\frac{1}{(n+1+z)^{\lambda+1}} \ll \frac{1}{(n+z)^{\lambda+1}} \quad \text{as } \lambda \to +\infty,
$$

for any $n = 0, 1, 2, \ldots$, the expansion in (27) describes completely the behavior of $G_0(\lambda, z)$ as $\lambda \to +\infty$. In particular (27) implies that for any fixed $z \in \mathbb{C}_0$ we have

$$
G_0(\lambda, z) = e^{\pi i \lambda} \frac{\Gamma(\lambda + 1)}{z^{\lambda + 1}} \left[1 + O\left(\left(\frac{z}{z + 1} \right)^{\lambda + 1} \right) \right], \quad \lambda \to +\infty.
$$
 (28)

3.2 Asymptotics of $G_1(\lambda, z)$ as $\lambda \to +\infty$

Let us, first, introduce some notation. We set

$$
\psi(t) := t \ln t, \qquad t \in [1, \infty)
$$
\n(29)

and

$$
\omega(\lambda) := \psi^{-1}(\lambda), \qquad \lambda \in [0, \infty). \tag{30}
$$

The function $\omega(\lambda)$ is well defined, smooth, and strictly increasing on $[0, \infty)$, with $\omega(0) = 1$, since $\psi(t)$ is smooth and strictly increasing on $[1, \infty)$, with $\psi(1) = 0$. Furthermore, it is easy to see that

$$
\omega(\lambda) = \frac{\lambda}{\ln \lambda} \left[1 + O\left(\frac{\ln \ln \lambda}{\ln \lambda}\right) \right], \qquad \lambda \to +\infty.
$$
 (31)

Lemma 1. For any fixed $z \in \mathbb{C}_0$ and for any fixed α such that

$$
0 < \alpha < \frac{1}{2},\tag{32}
$$

the function $G_1(\lambda, z)$ of (24) satisfies

$$
G_1(\lambda, z) = \frac{\sqrt{\pi} e^{\lambda \ln \ln \omega(\lambda)} e^{-\omega(\lambda)}}{\sqrt{A}} \omega(\lambda)^z \left[1 + O\left(\frac{1}{\lambda^{\alpha}}\right)\right], \qquad \lambda \to +\infty,
$$
 (33)

where $\omega(\lambda)$ is given by (29)-(30) and

$$
A := \lambda \frac{1 + \ln \omega(\lambda)}{2 \ln^2 \omega(\lambda)} \tag{34}
$$

(thus, from (31) we have that $A \sim \lambda/(2 \ln \lambda)$ as $\lambda \to +\infty$).

Proof. The integral of (24) can be viewed as a "Laplace integral" with large parameter λ [1]. Thus, we will try to apply the so-called Laplace method for integrals.

We begin by looking for the value of *t* which minimizes the function

$$
h(t) := t - \lambda \ln \ln t, \qquad t \in [1, \infty), \tag{35}
$$

appearing in the exponent of the integrand in (24). Since

$$
h'(t) = 1 - \frac{\lambda}{t \ln t} \tag{36}
$$

we have that $h(t)$ has a unique minimum attained when

$$
t\ln t = \lambda
$$

namely when (recall (29)-(30))

$$
t = \omega(\lambda) \tag{37}
$$

and, since $\omega(\lambda) > 1$ for $\lambda > 0$, this minimum is attained in the interior of the interval $[1, \infty)$, for every $\lambda > 0$.

From the theory of Laplace integrals we know that the main contribution to the value of the integral in (24), as λ gets large, comes from the values of *t* around $\omega(\lambda)$. In order to avoid the dependence in λ we make the substitution $t = \omega(\lambda)x$. Then, formula (24) becomes

$$
G_1(\lambda, z) = \omega(\lambda)^z \int_{1/\omega(\lambda)}^{\infty} e^{-p(x)} x^{z-1} dx,
$$
\n(38)

where for typographical convenience we have set

$$
p(x) := h(\omega(\lambda)x) = \omega(\lambda)x - \lambda \ln(\ln \omega(\lambda) + \ln x), \qquad x \in [1/\omega(\lambda), \infty). \tag{39}
$$

The minimum of $p(x)$ on $[1/\omega(\lambda), \infty)$ is attained at $x = 1$. Since $p'(1) = 0$ the Taylor expansion of $p(x)$ (with remainder) about $x = 1$ up to the cubic term is

$$
p(x) = p(1) + \frac{p''(1)}{2}(x-1)^2 + \frac{p'''(c)}{6}(x-1)^3,
$$
\n(40)

where *c* is some number between 1 and *x*. From (39) we have

$$
p(1) = \omega(\lambda) - \lambda \ln \ln \omega(\lambda), \qquad p''(1) = \lambda \frac{1 + \ln \omega(\lambda)}{\ln^2 \omega(\lambda)}, \qquad (41)
$$

$$
p'''(c) = -\lambda \left\{ \frac{2[\ln \omega(\lambda) + \ln c]^2 + 3[\ln \omega(\lambda) + \ln c] + 2}{c^3[\ln \omega(\lambda) + \ln c]^3} \right\}.
$$
 (42)

In particular, formula (42), combined with (31), imply that

$$
p'''(c) = O\left(\frac{\lambda}{\ln \lambda}\right) \quad \text{as } \lambda \to +\infty,
$$
 (43)

provided $\delta \leq x \leq 1/\delta$ for some fixed $\delta \in (0,1)$. Moreover, by (40) and (41) we can see that the quadratic Taylor polynomial of $p(x)$ about $x = 1$ is

$$
P(x; \lambda) := \omega(\lambda) - \lambda \ln \ln \omega(\lambda) + \lambda \frac{1 + \ln \omega(\lambda)}{2 \ln^2 \omega(\lambda)} (x - 1)^2
$$
 (44)

and it follows easily from (40), (41), (42), (44), and (32) that for all λ sufficiently large we must have

$$
|p(x) - P(x; \lambda)| \le \frac{1}{\lambda^{3\alpha - 1}}, \qquad x \in I_{\lambda}, \tag{45}
$$

where I_{λ} is the interval

$$
I_{\lambda} := \left[1 - \frac{1}{\lambda^{\alpha}}, 1 + \frac{1}{\lambda^{\alpha}}\right].
$$
 (46)

From now on let us assume, without loss of generality, that $\alpha > 1/3$, so that $1/\lambda^{3\alpha-1} \to 0$ as $\lambda \to +\infty$. Then, from (45) it follows that

$$
e^{-p(x)} = e^{-P(x;\lambda)} + R(x;\lambda), \qquad x \in I_{\lambda}, \tag{47}
$$

where

$$
|R(x;\lambda)| \le \frac{2}{\lambda^{3\alpha - 1}} e^{-P(x;\lambda)}, \qquad x \in I_{\lambda}, \tag{48}
$$

for all sufficiently large λ .

Multiplying by x^{z-1} and then taking integrals in (47) yields

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-p(x)} x^{z-1} dx = \int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-P(x;\lambda)} x^{z-1} dx + \int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} R(x;\lambda) x^{z-1} dx.
$$
\n(49)

Having formula (49), in order to complete the proof of the lemma, we need to verify the following three claims.

Claim 1. For the first integral in the right-hand side of (49) we have

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-P(x;\lambda)} x^{z-1} dx = \frac{\sqrt{\pi} e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)}}{\sqrt{A}} \left[1 + O\left(\frac{1}{\lambda^{\alpha}}\right)\right], \quad \lambda \to +\infty.
$$
\n(50)

and

Claim 2. For the second integral in the right-hand side of (49) we have

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} R(x;\lambda)x^{z-1}dx = \frac{e^{-\omega(\lambda)+\lambda\ln\ln\omega(\lambda)}}{\sqrt{A}} O\left(\frac{1}{\lambda^{3\alpha-1}}\right), \qquad \lambda \to +\infty. \tag{51}
$$

Claim 3. For the integral in the left-hand side of (49) we have

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-p(x)} x^{z-1} dx = \int_{1/\omega(\lambda)}^{\infty} e^{-p(x)} x^{z-1} dx + \frac{e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)}}{\sqrt{A}} o\left(\frac{1}{\lambda^{\alpha}}\right)
$$
(52)

as $\lambda \to +\infty$.

Proof of Claim 1. In view of (44) we have

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-P(x;\lambda)} x^{z-1} dx = e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)} \int_{-\lambda^{-\alpha}}^{\lambda^{-\alpha}} e^{-\lambda \frac{1+\ln \omega(\lambda)}{2\ln^2 \omega(\lambda)} \xi^2} (1+\xi)^{z-1} d\xi.
$$
\n(53)

Now, for $|\xi| \leq \lambda^{-\alpha}$ we have (since *z* is fixed)

$$
(1+\xi)^{z-1} = 1 + O\left(\frac{1}{\lambda^{\alpha}}\right), \qquad \lambda \to +\infty.
$$
 (54)

Thus, (53) implies

$$
\int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-P(x;\lambda)} x^{z-1} dx = e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)} \left[1 + O\left(\frac{1}{\lambda^{\alpha}}\right) \right] \int_{-\lambda^{-\alpha}}^{\lambda^{-\alpha}} e^{-A\xi^2} d\xi,
$$
\n(55)

where *A* is given by (34).

Finally, since from (32), (31), and (34) we have that $\sqrt{A} \lambda^{-\alpha} \gg \lambda^{(1-2\alpha)/4} \rightarrow$ +*∞*, while

$$
\int_0^{\sqrt{A}\,\lambda^{-\alpha}} e^{-u^2} du = \frac{\sqrt{\pi}}{2} + O\left(\frac{e^{-A\,\lambda^{-2\alpha}}}{\sqrt{A}\,\lambda^{-\alpha}}\right) \tag{56}
$$

as $\lambda \rightarrow +\infty$, formula (55) implies (50). *Proof of Claim 2*. Using (48) we get

$$
\left| \int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} R(x;\lambda) x^{z-1} dx \right| \leq \frac{2}{\lambda^{3\alpha-1}} \int_{1-\lambda^{-\alpha}}^{1+\lambda^{-\alpha}} e^{-P(x;\lambda)} x^{\Re(z)-1} dx,\qquad(57)
$$

thus (51) follows by using (50) in (57) .

Proof of Claim 3. We assume, without loss of generality, that $\lambda > 1$. Then

$$
\left| \int_{1+\lambda^{-\alpha}}^{\infty} e^{-p(x)} x^{z-1} dx \right| \leq \int_{1+\lambda^{-\alpha}}^{2} e^{-p(x)} x^{\Re(z)-1} dx + \int_{2}^{\infty} e^{-p(x)} x^{\Re(z)-1} dx \tag{58}
$$

From (39) we have that $p''(x) > 0$ for all $x \in [1/\omega(\lambda), \infty)$, while $p'(1) = 0$ by (29)-(30). Thus, $p(x)$ is convex on $[1/\omega(\lambda), \infty)$ and increasing on $[1, \infty)$. Consequently, for the first integral in the right-hand side of (58) we have

$$
\int_{1+\lambda^{-\alpha}}^{2} e^{-p(x)} x^{\Re(z)-1} dx \le 2^{\Re(z)} \int_{1+\lambda^{-\alpha}}^{2} e^{-p(x)} dx
$$

$$
\le 2^{\Re(z)} e^{-p(1+\lambda^{-\alpha})} \int_{1+\lambda^{-\alpha}}^{2} e^{-p'(1+\lambda^{-\alpha})(x-1-\lambda^{-\alpha})} dx
$$

(59)

(since $p(x) \ge p(1+\lambda^{-\alpha})+p'(1+\lambda^{-\alpha})(x-1-\lambda^{-\alpha})$ on $[1/\omega(\lambda), \infty)$ by convexity). Formula (59) implies

$$
\int_{1+\lambda^{-\alpha}}^{2} e^{-p(x)} x^{\Re(z)-1} dx \le 2^{\Re(z)} \frac{e^{-p(1+\lambda^{-\alpha})}}{p'(1+\lambda^{-\alpha})}.
$$
 (60)

By (45) we have

$$
p(1 + \lambda^{-\alpha}) = \omega(\lambda) - \lambda \ln \ln \omega(\lambda) + \frac{1 + \ln \omega(\lambda)}{2 \ln^2 \omega(\lambda)} \lambda^{1 - 2\alpha} + O\left(\frac{1}{\lambda^{3\alpha - 1}}\right), \quad (61)
$$

thus

$$
e^{-p(1+\lambda^{-\alpha})} = e^{-\omega(\lambda)+\lambda\ln\ln\omega(\lambda)} e^{-\frac{1+\ln\omega(\lambda)}{2\ln^2\omega(\lambda)}\lambda^{1-2\alpha}} \left[1+O\left(\frac{1}{\lambda^{3\alpha-1}}\right)\right].
$$
 (62)

Furthermore, from the Taylor expansion with remainder of $p'(x)$ about $x = 1$, namely

$$
p'(x) = p'(1) + p''(1)(x - 1) + \frac{p'''(c)}{2}(x - 1)^2, \qquad x \in I_\lambda,
$$

together with (41) , (43) , and the fact that $p'(1) = 0$, we obtain

$$
p'(1 + \lambda^{-\alpha}) = \frac{\lambda^{1-\alpha}}{\ln \omega(\lambda)} \left[1 + O\left(\lambda^{1-2\alpha}\right) \right].
$$
 (63)

Hence, by using (62) and (63) in (60) (and recalling (31) and the fact that $1 - 2\alpha > 0$) it follows that

$$
\int_{1+\lambda^{-\alpha}}^{2} e^{-p(x)} x^{\Re(z)-1} dx = \frac{e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)}}{\sqrt{A}} o\left(\frac{1}{\lambda^{\alpha}}\right), \qquad \lambda \to +\infty. \tag{64}
$$

As for the second integral in the right-hand side of (58), it is easy to see that it is much smaller than the bound given by (64) for the first integral in the right-hand side of (58). We can see that, e.g., by writing

$$
\int_{2}^{\infty} e^{-p(x)} x^{\Re(z) - 1} dx = \int_{2}^{\infty} e^{-[p(x) - p(x - 3/2)]} e^{-p(x - 3/2)} x^{\Re(z) - 1} dx \tag{65}
$$

and observing that the function $e^{-p(x-3/2)}x^{\Re(z)-1}$ is bounded on [2, ∞) uniformly in λ (for, say, $\lambda > 2$).

Finally, in the same manner we can show that

$$
\int_{1/\omega(\lambda)}^{1-\lambda^{-\alpha}} e^{-p(x)} x^{z-1} dx = \frac{e^{-\omega(\lambda)+\lambda \ln \ln \omega(\lambda)}}{\sqrt{A}} o\left(\frac{1}{\lambda^{\alpha}}\right), \qquad \lambda \to +\infty. \tag{66}
$$

Therefore, (52) follows immediately by using (64) and (66) in (58) .

If we use the estimate (31) in (33) we get the following corollary

Corollary 1. For the function $G_1(\lambda, z)$ of (24) we have

$$
G_1(\lambda, z) \sim \sqrt{\frac{\pi}{2}} \left[\ln \omega(\lambda) \right]^{\lambda} e^{-\omega(\lambda)} \left(\frac{\lambda}{\ln \lambda} \right)^{z - (1/2)} \quad \text{as } \lambda \to +\infty,
$$
 (67)

where $\omega(\lambda)$ is given by (29)-(30).

Remark 1. For the function $G_1(\lambda, z)$ of (24) we have

$$
G_1(m, z) = \Gamma^{(m)}(z, 1), \qquad m = 0, 1, \dots,
$$
\n(68)

where

$$
\Gamma(z,1) := \int_1^\infty e^{-t} t^{z-1} dt, \qquad z \in \mathbb{C},
$$
\n(69)

is the so-called upper incomplete Gamma function. Since $\Gamma(z, 1)$ is entire in z, the radius of convergence of its Taylor series about any z_0 is infinite, and this agrees with the asymptotics of $G_1(\lambda, z)$ for large λ , as given by (67). As for the function $G_0(\lambda, z)$ of (23) we have

$$
G_0(m, z) = \gamma^{(m)}(z, 1), \qquad m = 0, 1, \dots,
$$
\n(70)

where

$$
\gamma(z,1) := \int_0^1 e^{-t} t^{z-1} dt, \qquad \Re(z) > 0,
$$
\n(71)

is the so-called lower incomplete Gamma function. Since $\gamma(z, 1)$ has a simple pole at $z = 0$, the radius of convergence of its Taylor series about any z_0 with $\Re(z) > 0$ is $|z_0|$, and this agrees with the asymptotics of $G_0(\lambda, z)$ for large λ , as given by (28). Thus $G_0(\lambda, z)$ grows much faster than $G_1(\lambda, z)$ as $\lambda \to \infty$. This is, also, reflected in the corollary that follows.

Corollary 2. For the function $G(\lambda, z)$ of (3) we have

$$
G(\lambda, z) = e^{\pi i \lambda} \Gamma(\lambda + 1) \sum_{n=0}^{\infty} \frac{(-1)^n}{n! (n + z)^{\lambda + 1}}
$$

+ $\sqrt{\frac{\pi}{2}} \left[\ln \omega(\lambda) \right]^{\lambda} e^{-\omega(\lambda)} \left(\frac{\lambda}{\ln \lambda} \right)^{z - (1/2)} [1 + o(1)], \qquad \lambda \to +\infty,$ (72)

where *z* is fixed with $\Re(z) > 0$ and $\omega(\lambda)$ is given by (29)-(30).

Corollary 2 follows immediately from (27) and Corollary 1. Since $G(m, z)$ = $\Gamma^{(m)}(z)$ we obtain immediately from (72) the behavior of $\Gamma^{(m)}(z)$ as $m \to \infty$.

Corollary 3. For $\Re(z) > 0$ we have

$$
\Gamma^{(m)}(z) = (-1)^m m! \sum_{n=0}^{\infty} \frac{(-1)^n}{n! (n+z)^{m+1}} + \sqrt{\frac{\pi}{2}} \left[\ln \omega(m) \right]^m e^{-\omega(m)} \left(\frac{m}{\ln m} \right)^{z - (1/2)} [1 + o(1)], \qquad \lambda \to +\infty,
$$
\n(73)

where $\omega(\cdot)$ is given by (29)-(30). In particular,

$$
\Gamma^{(m)}(z) \sim \frac{(-1)^m m!}{z^{m+1}}, \qquad \lambda \to +\infty. \tag{74}
$$

3.3 Examples

1. If we set $z = 1$ in (73), we obtain

$$
\Gamma^{(m)}(1) = (-1)^m m! \sum_{n=0}^{\infty} \frac{(-1)^n}{n! (n+1)^{m+1}} + \sqrt{\frac{\pi}{2}} \left[\ln \omega(m) \right]^m e^{-\omega(m)} \sqrt{\frac{m}{\ln m}} \left[1 + o(1) \right], \qquad \lambda \to +\infty, \tag{75}
$$

in particular, $\Gamma^{(m)}(1) \sim (-1)^m m!$. Also, in view of (68), formula (67) gives

$$
\Gamma^{(m)}(1,1) \sim \sqrt{\frac{\pi}{2}} \left[\ln \omega(m) \right]^m e^{-\omega(m)} \sqrt{\frac{m}{\ln m}} \quad \text{as } m \to \infty,
$$
 (76)

where $\omega(\cdot)$ is given by (29)-(30).

2. Suppose we want the asymptotics of $G(\lambda, z)$ as $\lambda \to \infty$, in the case where z is a given complex number with $\Re(z) < 0$. Then, we can employ formula (19) or (20). For instance, for $z = -1/2$ formula (20) gives

$$
G\left(\lambda - 1, -\frac{1}{2}\right) = -2^{\lambda} \Gamma(\lambda) - 2^{\lambda} e^{-\pi i \lambda} \int_{-1/2}^{0} x^{\lambda - 1} G(\lambda, x + 1) dx,
$$

or

$$
G\left(\lambda - 1, -\frac{1}{2}\right) = -2^{\lambda} \Gamma(\lambda) - e^{-\pi i \lambda} \int_0^1 (\xi - 1)^{\lambda - 1} G\left(\lambda, \frac{\xi + 1}{2}\right) d\xi. \tag{77}
$$

We can now use (72) in (77) and get (as $\lambda \to +\infty$)

$$
G\left(\lambda - 1, -\frac{1}{2}\right) = -2^{\lambda}\Gamma(\lambda) + e^{\pi i \lambda} 2^{\lambda + 1} \Gamma(\lambda + 1) J(\lambda)
$$

$$
+ \left[\sqrt{\frac{\pi}{2}} + o(1)\right] \left[\ln \omega(\lambda)\right]^{\lambda} e^{-\omega(\lambda)} \int_0^1 (\xi - 1)^{\lambda - 1} \left(\frac{\lambda}{\ln \lambda}\right)^{\xi/2} d\xi
$$
(78)

where we have set

$$
J(\lambda) := \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \int_0^1 \frac{(1-\xi)^{\lambda-1}}{(2n+1+\xi)^{\lambda+1}} d\xi.
$$
 (79)

By invoking the formula

$$
-\frac{1}{(a+b)\lambda} \frac{d}{d\xi} \left[\frac{(a-\xi)^{\lambda}}{(b+\xi)^{\lambda}} \right] = \frac{(a-\xi)^{\lambda-1}}{(b+\xi)^{\lambda+1}},
$$
(80)

(79) becomes

$$
J(\lambda) = \frac{1}{2\lambda} \sum_{n=0}^{\infty} \frac{(-1)^n}{(n+1)!} \frac{1}{(2n+1)^{\lambda}}.
$$
 (81)

Also, it is not hard to show that

$$
\int_0^1 (\xi - 1)^{\lambda - 1} \left(\frac{\lambda}{\ln \lambda} \right)^{\xi/2} d\xi \sim \frac{1}{\lambda}, \qquad \lambda \to +\infty.
$$
 (82)

Using (81) and (82) in (78) yields

$$
G\left(\lambda - 1, -\frac{1}{2}\right) = \left(e^{\pi i \lambda} - 1\right) 2^{\lambda} \Gamma(\lambda) + e^{\pi i \lambda} 2^{\lambda} \Gamma(\lambda) \sum_{n=1}^{\infty} \frac{(-1)^n}{(n+1)!} \frac{1}{(2n+1)^{\lambda}}
$$

$$
+ \frac{\left[\ln \omega(\lambda)\right]^{\lambda}}{\lambda} e^{-\omega(\lambda)} \left[\sqrt{\frac{\pi}{2}} + o(1)\right], \qquad \lambda \to +\infty. \tag{83}
$$

By analytic continuation we have that $G(m, z) = \Gamma^{(m)}(z)$ for all $z \neq 0, -1, \ldots$. Hence, (83) implies

$$
\Gamma^{(m)}\left(-\frac{1}{2}\right) = 2^{m+1}m! \left\{ \left[(-1)^{m+1} - 1\right] - (-1)^m \sum_{n=1}^{\infty} \frac{(-1)^n}{(n+1)!} \frac{1}{(2n+1)^{m+1}} \right\} + \frac{\left[\ln \omega(m+1)\right]^{m+1}}{m+1} e^{-\omega(m+1)} \left[\sqrt{\frac{\pi}{2}} + o(1)\right], \qquad m \to \infty,
$$
\n(84)

in particular,

$$
\Gamma^{(2k)}\left(-\frac{1}{2}\right) \sim -2^{2k+2}(2k)!
$$
 and $\Gamma^{(2k+1)}\left(-\frac{1}{2}\right) \sim -\frac{2^{2k+1}(2k+1)!}{3^{2k+2}}$

as $k \to \infty$. The above asymptotic formulas are in agreement with the fact that for $z \in (-1,0)$ the "dominant" part of $\Gamma(z)$ is

$$
D_{\Gamma}(z) := \frac{1}{z} - \frac{1}{z+1}
$$

for which we have $D_{\Gamma}^{(2k)}$ $D_{\Gamma}^{(2k)}(-1/2) = -2^{2k+2}(2k)!$, while $D_{\Gamma}^{(2k+1)}$ $\Gamma^{(2\kappa+1)}(-1/2) = 0.$ Let us also mention that from the functional equation

$$
\Gamma^{(m)}(z+1) = z\Gamma^{(m)}(1) + m\Gamma^{(m-1)}(z), \qquad z \in \mathbb{C}, \quad m = 0, 1, ..., \tag{85}
$$

and the limits

$$
\Gamma^{(2k)}(0^-) = \Gamma^{(2k)}(-1^+) = -\infty, \qquad \Gamma^{(2k+1)}(0^-) = -\Gamma^{(2k)}(-1^+) = -\infty, \tag{86}
$$

it follows that $\Gamma^{(2k)}(x) < 0$ for $x \in (-1,0)$ and hence that $\Gamma^{(2k+1)}(x)$ is decreasing on $(-1,0)$. Consequently, for each $k = 0,1,...$ the function $\Gamma^{(2k+1)}(z)$ has a unique zero, say η_k , in (-1,0). A natural question here is whether η_k , $k \geq 0$, is a monotone sequence and, furthermore, whether $\eta_k \to -1/2$ as $k \to \infty$.

3. Let us, finally, consider the case where *z* is real and positive. Then formula (2) implies that $\Gamma^{(2k)}(z) > 0$ for all $k \geq 0$. Consequently, all odd derivatives $\Gamma^{(2k+1)}(z)$, $k \geq 0$, are increasing in $(0, \infty)$, with $\Gamma^{(2k+1)}(0^+) = -\infty$ and $\Gamma^{(2k+1)}(+\infty) = +\infty$. Hence, for each $k \geq 0$ there is a unique $\zeta_k \in (0, \infty)$ such that

$$
\Gamma^{(2k+1)}(\zeta_k) = 0.\tag{87}
$$

Suppose that the sequence ζ_k , $k \geq 0$, has a bounded subsequence, namely there is an $M > 0$ such that $\zeta_k < M$ for infinitely many values of k . Then, we should have $\Gamma^{(2k+1)}(M) > 0$ for infinitely many values of k. But this is impossible, since formula (73), or just (74), implies that $\Gamma^{(2k+1)}(M) < 0$ for all sufficiently large *k*. Therefore $\zeta_k \to \infty$. A natural question here is whether ζ_k is increasing.

References

[1] C.M. Bender and S.A. Orszag, *Advanced Mathematical Methods for Scientists and Engineers I: Asymptotic Methods and Perturbation Theory*, Springer-Verlag, New York, 1999.